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Computational Fluid Dynamics

Finite Difference Method

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FDM: Basics

- The finite difference method is based on the properties of the Taylor series and on the straight forward application of the definition of derivatives.
- The earliest work in this field is due to Euler in 1768.
- For a function $u(x)$, the derivative at point x is defined as

$$u_x \equiv \frac{\partial u}{\partial x} = \lim_{\Delta x \rightarrow 0} \frac{u(x + \Delta x) - u(x)}{\Delta x}$$

- This means that
(Actual derivative) = (approximate derivative) + (truncation error)
- The power Δx with which the error goes to zero for $\Delta x \rightarrow 0$ is called the **order of the difference equation**.
- Expanding $u(x + \Delta x)$ around $u(x)$, we get

$$u(x + \Delta x) = u(x) + \Delta x \cdot u_x + \frac{(\Delta x)^2}{2!} u_{xx} + \dots \rightarrow$$
$$\frac{u(x + \Delta x) - u(x)}{\Delta x} = u_x + \frac{(\Delta x)}{2!} u_{xx} + \dots = u_x + O(\Delta x)$$

Various Differences

- There are three difference forms:

$$(u_x)_i = \frac{u_{i+1} - u_i}{\Delta x} + O(\Delta x) \quad \textit{forward}$$

$$(u_x)_i = \frac{u_i - u_{i-1}}{\Delta x} + O(\Delta x) \quad \textit{backward}$$

$$\left\{ \begin{array}{l} (u_x)_i = \frac{u_{i+1} - u_{i-1}}{2\Delta x} + O(\Delta x^2) \\ (u_x)_{i+1/2} = \frac{u_{i+1} - u_i}{\Delta x} + O(\Delta x^2) \\ (u_x)_{i-1/2} = \frac{u_i - u_{i-1}}{\Delta x} + O(\Delta x^2) \end{array} \right. \quad \textit{central}$$

Using Arbitrary Number of Points

- We can involve as many points as we want to obtain a required accuracy. For example, we can write:

$$(u_x)_i = \frac{au_i + bu_{i-1} + cu_{i-2}}{\Delta x} + O(\Delta x^2)$$

- The coefficients a , b and c are obtained as

$$u_{i-2} = u_i - 2\Delta x(u_x)_i + 2(\Delta x)^2(u_{xx})_i - \frac{(2\Delta x)^3}{6}(u_{xxx})_i + \dots$$

$$u_{i-1} = u_i - \Delta x(u_x)_i + \frac{(\Delta x)^2}{2}(u_{xx})_i - \frac{(\Delta x)^3}{6}(u_{xxx})_i + \dots$$

- Then,

$$au_i + bu_{i-1} + cu_{i-2} = (a+b+c)u_i - \Delta x(2c+b)(u_x)_i + \frac{(\Delta x)^3}{2}(4c+b)(u_{xx})_i + O(\Delta x^3)$$

- This requires

$$\begin{cases} a+b+c=0 \\ 2c+b=-1 \\ 4c+b=0 \end{cases} \rightarrow \begin{cases} a=3/2 \\ b=-2 \\ c=1/2 \end{cases} \rightarrow (u_x)_i = \frac{3u_i - 4u_{i-1} + u_{i-2}}{2\Delta x}$$

Using Arbitrary...

- A similar relation can be written using $i, i+1$ and $i+2$

$$(u_x)_i = \frac{-3u_i + 4u_{i+1} - u_{i+2}}{\Delta x} + O(\Delta x^2)$$

- For the second order derivatives, we can write:

$$(u_{xx})_i = \frac{(u_x)_{i+1} - (u_x)_i}{\Delta x} = \frac{u_{i+1} - 2u_i + u_{i-1}}{(\Delta x)^2} + O(\Delta x^2)$$

- Another form which is first order can be written as

$$(u_{xx})_i = \frac{u_i - 2u_{i-1} + u_{i-2}}{(\Delta x)^2} + O(\Delta x)$$

General Method for Deriving Formulas

- We use the following operator

$$E u_i = u_{i+1} \quad (\text{displacement operator})$$

$$E^{-1} u_i = u_{i-1}$$

$$E^n u_i = u_{i+n} \quad (\text{in general})$$

- Also use

$$\delta^+ : \text{forward difference operator} \quad \delta^+ u_i = u_{i+1} - u_i$$

$$\delta^- : \text{backward difference operator} \quad \delta^- u_i = u_i - u_{i-1}$$

$$\delta : \text{central difference operator} \quad \delta u_i = u_{i+1/2} - u_{i-1/2}$$

$$\bar{\delta} : \text{central difference operator} \quad \bar{\delta} u_i = (u_{i+1} - u_{i-1})/2$$

$$\mu : \text{averaging operator} \quad \mu u_i = (u_{i+1/2} + u_{i-1/2})/2$$

$$D : \text{differential operator} \quad D u_i = \frac{\partial u}{\partial x} = u_x$$

- We note that

$$\delta^+ = E - 1$$

$$\delta^- = 1 - 1/E$$

General Method...

- This results in the following relations:

$$\delta^- = E^{-1} \delta^+$$

$$\delta^+ \delta^- = \delta^- \delta^+ = \delta^+ - \delta^- = \delta^2$$

$$\delta = E^{1/2} - E^{-1/2} = (E - E^{-1})/2$$

$$\mu = (E^{1/2} + E^{-1/2})/2$$

- For n repeated actions of the operator δ , we have

$$(\delta^+)^2 = \delta^+ \delta^+ = E^2 - 2E + 1$$

$$(\delta^+)^3 = (E - 1)^3 = E^3 - 3E^2 + 3E - 1$$

Difference Formulas for First Derivatives

- Using the Taylor expansion, we get

$$u(x + \Delta x) = u(x) + \Delta x \cdot u_x + \frac{(\Delta x)^2}{2!} u_{xx} + \frac{(\Delta x)^3}{3!} u_{xxx} + \dots$$

- Or

$$E u(x) = \left(1 + \Delta x \cdot D + \frac{(\Delta x D)^2}{2!} + \frac{(\Delta x D)^3}{3!} + \dots \right) u(x) \quad \rightarrow$$
$$E u(x) = e^{\Delta x D} u(x) \quad \rightarrow \quad E = \exp(\Delta x D) \quad \rightarrow \quad \Delta x D = \ln(E)$$

- Forward Difference:** from the above relations, we can write

$$\Delta x D = \ln(E) = \ln(1 + \delta^+) = \delta^+ - \frac{(\delta^+)^2}{2} + \frac{(\delta^+)^3}{3} - \frac{(\delta^+)^4}{4} + \dots$$

- Keeping the first three terms, we get

$$(u_x)_i = D u_i = \frac{-3u_i + 4u_{i+1} - u_{i+2}}{2 \Delta x} + \frac{(\Delta x)^2}{3} u_{xxx}$$

Difference formulas...

- **Backward difference:**

$$\Delta x D = \ln(E) = -\ln(1 - \delta^-) = \delta^- + \frac{(\delta^-)^2}{2} + \frac{(\delta^-)^3}{3} + \frac{(\delta^-)^4}{4} + \dots$$

or

$$(u_x)_i = D u_i = \frac{3u_i - 4u_{i-1} + u_{i-2}}{2\Delta x} + \frac{(\Delta x)^2}{3} u_{xxx}$$

- **Central Difference:**

$$\begin{aligned} \delta u_i &= u_{i+1/2} - u_{i-1/2} = (E^{1/2} - E^{-1/2})u_i \rightarrow \\ \delta &= \exp(\Delta x D / 2) - \exp(-\Delta x D / 2) = 2 \sinh(\Delta x D / 2) \rightarrow \\ \Delta x D &= 2 \sinh^{-1}(\delta / 2) = 2 \left[\frac{\delta}{2} - \frac{1}{2 \times 3} \left(\frac{\delta}{2}\right)^3 + \frac{1 \times 3}{2 \times 4 \times 5} \left(\frac{\delta}{2}\right)^5 - \frac{1 \times 3 \times 5}{2 \times 4 \times 6 \times 7} \left(\frac{\delta}{2}\right)^7 + \dots \right] \\ &= \delta - \frac{\delta^3}{24} + \frac{3\delta^5}{640} - \frac{5\delta^7}{7168} + \dots \end{aligned}$$

- **Keeping the first term, we obtain**

$$(u_x)_i = \frac{u_{i+1/2} - u_{i-1/2}}{\Delta x} - \frac{(\Delta x)^2}{24} u_{xxx}$$

Difference formulas...

- Keeping the first two terms, we obtain:

$$(u_x)_i = \frac{-u_{i+3/2} + 27u_{i+1/2} - 27u_{i-1/2} + u_{i-3/2}}{24\Delta x} + \frac{3}{640}(\Delta x)^4 \frac{\partial^5 u}{\partial x^5}$$

- We could use $\bar{\delta}$ to get

$$\bar{\delta} = \frac{1}{2}(E - E^{-1}) = \frac{1}{2}(e^{\Delta x D} - e^{-\Delta x D}) = \sinh(\Delta x D) \rightarrow$$

$$\Delta x D = \sinh^{-1}(\bar{\delta}) = \bar{\delta} - \frac{(\bar{\delta})^3}{6} + \frac{3}{2 \times 4 \times 5}(\bar{\delta})^5 - \dots$$

- Nothing that $\mu^2 = 1 + \delta^2 / 4$ or

$$1 = \mu(1 + \delta^2 / 4)^{-1/2} = \mu \left(1 - \frac{\delta^2}{8} + \frac{3\delta^4}{128} - \frac{5\delta^6}{1024} + \dots \right)$$

- We obtain

$$\Delta x D = \mu \left(\delta - \frac{1}{3!} \delta^3 + \frac{1^2 \times 2^2}{5!} \delta^5 - \dots \right) = \bar{\delta} \left(1 - \frac{1}{3!} \delta^2 + \frac{2^2}{5!} \delta^4 - \frac{2^2 \times 3^2}{7!} \delta^6 + \dots \right)$$

Difference formulas...

- Then, keeping one term, we get

$$(u_x)_i = \frac{u_{i+1} - u_{i-1}}{2 \Delta x} - \frac{(\Delta x)^2}{6} u_{xxx}$$

- And keeping two terms, we get

$$(u_x)_i = \frac{-u_{i+2} + 8u_{i+1} - 8u_{i-1} + u_{i-2}}{12 \Delta x} + \frac{(\Delta x)^4}{30} u_{xxxx}$$

Higher Order Derivatives

- For higher-order derivatives, we can use a one-sided forward difference as

$$\begin{aligned}\left(\frac{\partial^n u}{\partial x^n}\right)_i &= D^n u_i = \frac{1}{\Delta x^n} [\ln(1 + \delta^+)]^n u_i \\ &= \frac{1}{\Delta x^n} \left[(\delta^+)^n - \frac{n}{2} (\delta^+)^{n+1} + \frac{n(3n+5)}{24} (\delta^+)^{n+2} - \frac{n(n+2)(n+3)}{48} (\delta^+)^{n+3} + \dots \right] u_i\end{aligned}$$

- This can be written in terms of the backward and central differences as

$$\left(\frac{\partial^n u}{\partial x^n}\right)_i = \frac{-1}{\Delta x^n} [\ln(1 - \delta^-)]^n u_i \quad \left(\frac{\partial^n u}{\partial x^n}\right)_i = \left[\frac{2}{\Delta x} \sinh^{-1}\left(\frac{\delta}{2}\right) \right]^n u_i$$

- For even values of n , the central difference form generates formulas using the integer mesh points. For odd values of n , the formulas involve points at half-integer mesh points.
- For the following equation, the inverse is true:

$$D^n u_i = \frac{\mu \delta^n}{\Delta x^n} \left(1 - \frac{n+3}{24} \delta^2 + \frac{5n^2 + 52n + 135}{5760} \delta^4 + \dots \right) u_i$$

Higher Order Derivatives

- For example, we obtain the followings:

$$(u_{xx})_i = \frac{1}{\Delta x^2} \left[(\delta^+)^2 - (\delta^+)^3 + \frac{11}{12} (\delta^+)^4 - \frac{5}{6} (\delta^+)^5 + \dots \right] u_i$$

$$(u_{xx})_i = \frac{1}{\Delta x^2} \left[(\delta^-)^2 + (\delta^-)^3 + \frac{11}{12} (\delta^-)^4 + \frac{5}{6} (\delta^-)^5 + \dots \right] u_i$$

$$(u_{xx})_i = \frac{1}{\Delta x^2} \left[(\delta)^2 - \frac{1}{12} (\delta)^4 + \frac{1}{90} (\delta)^6 - \frac{1}{560} (\delta)^8 + O(\Delta x^8) \right] u_i$$

- By keeping only the first term, these formulas give:

$$(u_{xx})_i = \frac{1}{\Delta x^2} (u_{i+2} - 2u_{i+1} + u_i) - \Delta x \cdot u_{xxx} \quad (1st \text{ order forward})$$

$$(u_{xx})_i = \frac{1}{\Delta x^2} (u_i - 2u_{i-1} + u_{i-2}) + \Delta x \cdot u_{xxx} \quad (1st \text{ order backward})$$

$$(u_{xx})_i = \frac{1}{\Delta x^2} (u_{i+1} - 2u_i + u_{i-1}) - \frac{\Delta x^2}{12} \cdot u_{xxxx} \quad (2nd \text{ order central})$$

$$(u_{xx})_i = \frac{1}{2\Delta x^2} (u_{i+3/2} - u_{i+1/2} - u_{i-1/2} + u_{i-3/2}) - \frac{5\Delta x^2}{24} \cdot u_{xxxx} \quad (2nd \text{ order central})$$

Higher Order Derivatives...

- The one-sided difference formulas are only first-order accurate, while the central difference always leads to a higher-order of accuracy.
- By keeping the first two terms of the above formulas, we obtain higher-order formulas:

$$(u_{xx})_i = \frac{1}{\Delta x^2}(2u_i - 5u_{i+1} + 4u_{i+2} - u_{i+3}) + \frac{11\Delta x^2}{12} \cdot u_{xxxx} \quad (2nd \text{ order forward})$$

$$(u_{xx})_i = \frac{1}{\Delta x^2}(2u_i - 5u_{i-1} + 4u_{i-2} - u_{i-3}) - \frac{11\Delta x^2}{12} \cdot u_{xxxx} \quad (2nd \text{ order backward})$$

$$(u_{xx})_i = \frac{1}{12\Delta x^2}(-u_{i+2} + 16u_{i+1} - 30u_i + 16u_{i-1} - u_{i-2}) + \frac{\Delta x^4}{90} \cdot u_{xxxxxx} \quad (4th \text{ order central})$$

$$(u_{xx})_i = \frac{1}{48\Delta x^2}(-5u_{i+5/2} + 39u_{i+3/2} - 34u_{i+1/2} - 34u_{i-1/2} + 39u_{i-3/2} - 5u_{i-5/2}) + \frac{259\Delta x^4}{5760} \cdot u_{xxxxxx} \quad (4th \text{ order central})$$

- For a term like $\frac{\partial}{\partial x} \left(k(x) \frac{\partial u}{\partial x} \right)$, we can use: (2nd order central)

$$\begin{aligned} \frac{\partial}{\partial x} \left[k(x) \frac{\partial u}{\partial x} \right] u_i &= \frac{1}{(\Delta x)^2} \delta^+ (k_{i-1/2} \delta^-) u_i + O(\Delta x^2) \\ &= \frac{k_{i+1/2}(u_{i+1} - u_i)}{(\Delta x)^2} - \frac{k_{i-1/2}(u_i - u_{i-1})}{(\Delta x)^2} + O(\Delta x^2) \end{aligned}$$

3rd Order Derivatives

- As an special case, we can easily show that

$$\begin{aligned}\left(\frac{\partial^3 u}{\partial x^3}\right)_i &= \frac{1}{(\Delta x)^3} (u_{i+3} - 3u_{i+2} + 3u_{i+1} - u_i) - \frac{\Delta x}{2} \frac{\partial^4 u}{\partial x^4} && \text{1st order forward} \\ \left(\frac{\partial^3 u}{\partial x^3}\right)_i &= \frac{1}{2(\Delta x)^3} (-3u_{i+4} + 14u_{i+3} - 24u_{i+2} + 18u_{i+1} - 5u_i) + \frac{21(\Delta x)^2}{12} \frac{\partial^5 u}{\partial x^5} && \text{2nd order forward} \\ \left(\frac{\partial^3 u}{\partial x^3}\right)_i &= \frac{1}{(\Delta x)^3} (u_i - 3u_{i-1} + 3u_{i-2} - u_{i-3}) + \frac{\Delta x}{2} \frac{\partial^4 u}{\partial x^4} && \text{1st order backward} \\ \left(\frac{\partial^3 u}{\partial x^3}\right)_i &= \frac{1}{2(\Delta x)^3} (5u_i - 18u_{i-1} + 24u_{i-2} - 14u_{i-3} + 3u_{i-4}) - \frac{21(\Delta x)^2}{12} \frac{\partial^5 u}{\partial x^5} && \text{2nd order backward} \\ \left(\frac{\partial^3 u}{\partial x^3}\right)_i &= \frac{1}{(\Delta x)^3} (u_{i+3/2} - 3u_{i+1/2} + 3u_{i-1/2} - u_{i-3/2}) - \frac{(\Delta x)^2}{8} \frac{\partial^5 u}{\partial x^5} && \text{2nd order central} \\ \left(\frac{\partial^3 u}{\partial x^3}\right)_i &= \frac{1}{8(\Delta x)^3} (-u_{i+5/2} + 13u_{i+3/2} - 34u_{i+1/2} + 34u_{i-1/2} - 13u_{i-3/2} + u_{i-5/2}) - \frac{37(\Delta x)^4}{1920} \frac{\partial^7 u}{\partial x^7} && \text{4th order central}\end{aligned}$$

- Also, using integer points, we get:

$$\begin{aligned}\left(\frac{\partial^3 u}{\partial x^3}\right)_i &= \frac{1}{2(\Delta x)^3} (u_{i+2} - 2u_{i+1} + 2u_{i-1} + u_{i-2}) - \frac{(\Delta x)^2}{4} \frac{\partial^5 u}{\partial x^5} && \text{2nd order central} \\ \left(\frac{\partial^3 u}{\partial x^3}\right)_i &= \frac{1}{8(\Delta x)^3} (-u_{i+3} + 8u_{i+2} - 13u_{i+1} + 13u_{i-1} - 8u_{i-2} + u_{i-3}) - \frac{7(\Delta x)^4}{120} \frac{\partial^7 u}{\partial x^7} && \text{4th order central}\end{aligned}$$

4th Order Derivatives

- Similarly, for the 4th order formulas, we get

$$\begin{aligned}\left(\frac{\partial^4 u}{\partial x^4}\right)_i &= \frac{1}{(\Delta x)^4}(u_{i+4} - 4u_{i+3} + 6u_{i+2} - 4u_{i+1} + u_i) - 2\Delta x \frac{\partial^5 u}{\partial x^5} && \text{1st order forward} \\ \left(\frac{\partial^4 u}{\partial x^4}\right)_i &= \frac{1}{(\Delta x)^4}(u_i - 4u_{i-1} + 6u_{i-2} - 4u_{i-3} + u_{i-4}) + 2\Delta x \frac{\partial^5 u}{\partial x^5} && \text{1st order backward} \\ \left(\frac{\partial^4 u}{\partial x^4}\right)_i &= \frac{1}{(\Delta x)^4}(u_{i+2} - 4u_{i+1} + 6u_i - 4u_{i-1} + u_{i-2}) - \frac{(\Delta x)^2}{6} \frac{\partial^6 u}{\partial x^6} && \text{2nd order central}\end{aligned}$$

Implicit FD Formulas

- Implicit formulas are defined as expressions where derivatives at different mesh points appear simultaneously.
- We can generate implicit formulas from previous explicit formulas.
- For example starting with

$$\begin{aligned}\Delta x D &= \mu \delta (1 + \delta^2 / 6)^{-1} + O(\Delta x^5) \rightarrow \\ (1 + \delta^2 / 6) D &= \frac{\mu \delta}{\Delta x} + O(\Delta x^4) \rightarrow \\ \frac{1}{6} [(u_x)_{i+1} + 4(u_x)_i + (u_x)_{i-1}] &= \frac{u_{i+1} - u_{i-1}}{2\Delta x} + O(\Delta x^4)\end{aligned}$$

- The last equation is called **rational fraction** or **Pade's differencing**.

Implicit FD Formulas...

- The use of this formulas results in the following tri-diagonal system:

$$\begin{bmatrix} \ddots & & & & 0 \\ 1 & 4 & 1 & & \\ & 1 & 4 & 1 & \\ & & 1 & 4 & 1 \\ 0 & & & \ddots & \end{bmatrix} \begin{bmatrix} 0 \\ \vdots \\ (u_x)_{i-1} \\ (u_x)_i \\ (u_x)_{i+1} \\ \vdots \\ 0 \end{bmatrix} = \frac{3}{\Delta x} \begin{bmatrix} 0 \\ \vdots \\ u_i - u_{i-2} \\ u_{i+1} - u_{i-1} \\ u_{i+2} - u_i \\ \vdots \\ 0 \end{bmatrix}$$

- The reason why the formulas are higher order is that each u_x depends on all values of u^i 's .

Derivation of Implicit Formulas

- If we use a three point expression, we can write a general form as

$$\begin{aligned} & a^+ u_{i+1} + a^0 u_i + a^- u_{i-1} + \\ & b^+ (u_x)_{i+1} + b^0 (u_x)_i + b^- (u_x)_{i-1} + \\ & c^+ (u_{xx})_{i+1} + c^0 (u_{xx})_i + c^- (u_{xx})_{i-1} = 0 \end{aligned}$$

- Developing all the variables in a Taylor series about point i , We get
- In order to get a 2nd order accuracy for the second derivative, we request the coefficients up to the third order derivatives of the truncation error to vanish, i.e.:

$$\begin{aligned} u_{i\pm 1} = & u_i \pm \Delta x (u_x)_i + \frac{(\Delta x)^2}{2} (u_{xx})_i \pm \frac{(\Delta x)^3}{6} (u_{xxx})_i + \frac{(\Delta x)^4}{24} (u_{xxxx})_i \\ & \pm \frac{(\Delta x)^5}{5!} (u_{xxxxx})_i + \frac{(\Delta x)^6}{6!} (u_{xxxxxx})_i \pm \dots \end{aligned}$$

Derivation of Implicit Formulas...

- For the first derivative, we get

$$\begin{aligned}(u_x)_{i\pm 1} &= (u_x)_i \pm \Delta x (u_{xx})_i + \frac{(\Delta x)^2}{2} (u_{xxx})_i \pm \frac{(\Delta x)^3}{6} (u_{xxxx})_i \\ &\quad + \frac{(\Delta x)^4}{24} (u_{xxxxx})_i \pm \frac{(\Delta x)^5}{5!} (u_{xxxxxx})_i + \frac{(\Delta x)^6}{6!} (u_{xxxxxxx})_i \pm \dots\end{aligned}$$

- And for the second derivative, we get

$$\begin{aligned}(u_{xx})_{i\pm 1} &= (u_{xx})_i \pm \Delta x (u_{xxx})_i + \frac{(\Delta x)^2}{2} (u_{xxxx})_i \pm \frac{(\Delta x)^3}{6} (u_{xxxxx})_i \\ &\quad + \frac{(\Delta x)^4}{24} (u_{xxxxxx})_i \pm \frac{(\Delta x)^5}{5!} (u_{xxxxxxx})_i + \frac{(\Delta x)^6}{6!} (u_{xxxxxxx})_i \pm \dots\end{aligned}$$

Derivation of Implicit Formulas...

- In order to get a 2nd order accuracy for the 2nd derivative, we request that the coefficients up to the 3rd order derivatives of the truncation error to vanish:

$$\left\{ \begin{array}{l} a^+ + a^0 - a^- = 0 \\ \Delta x(a^+ - a^-) + b^+ + b^0 + b^- = 0 \\ \frac{\Delta x^2}{2}(a^+ + a^-) + \Delta x(b^+ - b^-) + c^+ + c^0 + c^- = 0 \rightarrow \\ \frac{\Delta x^3}{6}(a^+ - a^-) + \frac{\Delta x^2}{2}(b^+ - b^-) + \Delta x(c^+ - c^-) = 0 \\ \left\{ \begin{array}{l} a^+ = \frac{1}{2\Delta x} \left[-5b^+ - b^- + \frac{2}{\Delta x}(2c^- - 4c^+ - c^0) \right] \\ a^0 = \frac{2}{\Delta x} \left[b^+ - b^- + \frac{1}{\Delta x}(c^+ + c^0 - c^-) \right] \\ a^- = \frac{1}{2\Delta x} \left[b^+ + 5b^- + \frac{2}{\Delta x}(2c^+ - 4c^- - c^0) \right] \\ b^0 = 2(b^+ + b^-) + \frac{6}{\Delta x}(c^+ - c^-) \end{array} \right. \end{array} \right.$$

Derivation of Implicit Formulas...

- And the truncation error R reduces to

$$\begin{aligned} R = & \frac{\Delta x^3}{24} \left[2(b^+ - b^-) + \frac{10}{\Delta x}(c^+ - c^-) - \frac{2}{\Delta x}c^0 \right] \frac{\partial^4 u}{\partial x^4} + \\ & \frac{\Delta x^4}{120} \left[2(b^+ - b^-) + \frac{14}{\Delta x}(c^+ - c^-) \right] \frac{\partial^5 u}{\partial x^5} + \\ & \frac{\Delta x^5}{6!} \left[4(b^+ - b^-) + \frac{28}{\Delta x}(c^+ + c^-) - \frac{2}{\Delta x}c^0 \right] \frac{\partial^6 u}{\partial x^6} + \\ & \frac{\Delta x^6}{7!} \left[4(b^+ + b^-) + \frac{36}{\Delta x}(c^+ - c^-) \right] \frac{\partial^7 u}{\partial x^7} + \\ & \frac{\Delta x^7}{8!} \left[6(b^+ - b^-) + \frac{54}{\Delta x}(c^+ + c^-) - \frac{2}{\Delta x}c^0 \right] \frac{\partial^8 u}{\partial x^8} \end{aligned}$$

- This gives a four parameter family of implicit relations.
- One parameter may be set arbitrarily to one since our original form is homogeneous.

Derivation of Implicit Formulas...

- If we choose to make the first three terms in R equal to zero, we obtain:

$$\begin{aligned}b^+ &= \frac{1}{\Delta x} (8c^+ + c^-) \\b^- &= \frac{1}{\Delta x} (c^+ + 8c^-) \\c^0 &= -4(c^+ + c^-)\end{aligned}$$

- Then, we get:

$$\begin{aligned}&\frac{3}{2\Delta x^2} (13 + 3\alpha)u_{i+1} - \frac{24}{\Delta x^2} (1 + \alpha)u_i + \frac{3}{2\Delta x^2} (3 + 13\alpha)u_{i-1} - \\&\frac{1}{\Delta x} (8 + \alpha)(u_x)_{i+1} - \frac{8}{\Delta x} (1 - \alpha)(u_x)_i + \frac{1}{\Delta x} (1 + 8\alpha)(u_x)_{i-1} + \\&(u_{xx})_{i+1} - 4(1 + \alpha)(u_{xx})_i + \alpha(u_{xx})_{i-1} = 0\end{aligned}$$

- Where $\alpha = c^- / c^+$ and

$$R = \frac{8\Delta x^5}{7!} (1 - \alpha) \frac{\partial^7 u}{\partial x^7} + \frac{\Delta x^6}{8!} (1 + \alpha) \frac{\partial^8 u}{\partial x^8}$$

Derivation of Implicit Formulas...

- The unique, implicit relation of order six is obtained from $\alpha = 1$:

$$\frac{24}{\Delta x^2}(u_{i+1} - 2u_i + u_{i-1}) - \frac{9}{\Delta x}[(u_x)_{i+1} - (u_x)_{i-1}] + (u_{xx})_{i+1} - 8(u_{xx})_i + (u_{xx})_{i-1} = 0$$

with

$$R = \frac{2}{8!} \Delta x^6 \frac{\partial^8 u}{\partial x^8}$$

- The unique 4th order relation is obtained from $c^+ = c^0 = c^- = 0$
and $\beta = b^- / b^+$:

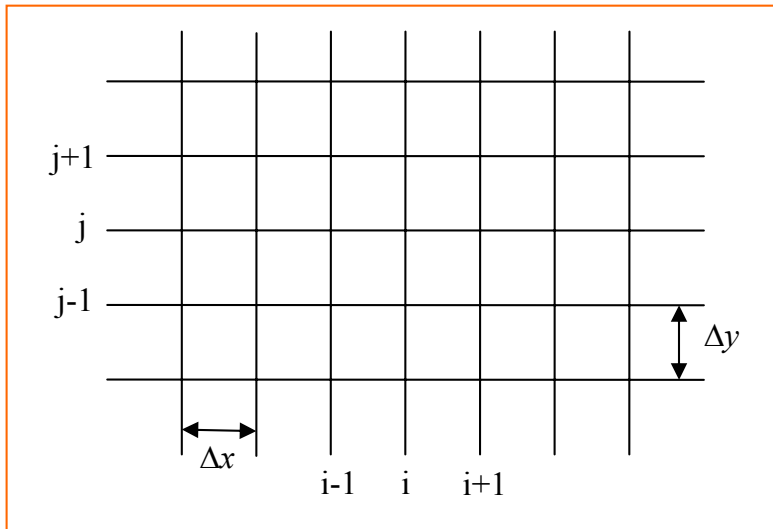
$$\frac{1}{2\Delta x}(-5 - \beta)u_{i+1} + \frac{2}{\Delta x}(1 - \beta)u_i + \frac{1}{2\Delta x}(1 + 5\beta)u_{i-1} + (u_x)_{i+1} + 2(1 + \beta)(u_x)_i + \beta(u_x)_{i-1} = 0$$

with

$$R = \frac{\Delta x^3}{12}(1 - \beta)\frac{\partial^4 u}{\partial x^4} + \frac{\Delta x^4}{60}(1 + \beta)\frac{\partial^5 u}{\partial x^5}$$

Note: SPECIAL CASE $\beta = 1$

Multi-dimensional FDM Formulas



- In a 2D space a rectangular mesh can be obtained by

$$x_i = x_0 + \Delta x \cdot i$$

$$y_i = y_0 + \Delta y \cdot j$$

Multi-dimensional FDM Formulas

- If we denote $u(x_i, y_j)$ by u_{ij} , we can apply the previous formulas on either variables x and y .
- For example: (1st order formulas)

$$\begin{aligned}(u_x)_{ij} &= \frac{u_{i+1,j} - u_{i,j}}{\Delta x} = \frac{1}{\Delta x} \delta_x^+ u_{ij} + O(\Delta x) \\ (u_y)_{ij} &= \frac{u_{i,j+1} - u_{i,j}}{\Delta y} = \frac{1}{\Delta y} \delta_y^+ u_{ij} + O(\Delta y)\end{aligned}$$

- And a second order formula can be written as

$$(u_{xx})_{ij} = \frac{u_{i+1,j} - 2u_{i,j} + u_{i-1,j}}{\Delta x^2} - \frac{\Delta x^2}{12} u_{xxxx}$$

Multi-dimensional FDM Formulas

- **Laplace Equation:**

$$(u_{xx} + u_{yy})_{ij} = \frac{u_{i-1,j} - 2u_{i,j} + u_{i+1,j}}{\Delta x^2} + \frac{u_{i,j-1} - 2u_{i,j} + u_{i,j+1}}{\Delta y^2} + O(\Delta x^2, \Delta y^2)$$

- **For $\Delta x = \Delta y$, we get**

$$\Delta u_{ij} = (u_{xx} + u_{yy})_{ij} = \frac{u_{i+1,j} + u_{i-1,j} + u_{i,j-1} + u_{i,j+1} - 4u_{ij}}{\Delta x^2} - \frac{\Delta x^2}{12} \left(\frac{\partial^4 u}{\partial x^4} + \frac{\partial^4 u}{\partial y^4} \right)$$

- **A more general term $\nabla \cdot (k \nabla)$ can be written as**

$$\nabla \cdot (k \nabla) = \frac{1}{\Delta x^2} (\delta_x^+ k_{i+1/2,j} \delta_x^-) u_{ij} + \frac{1}{\Delta y^2} (\delta_y^+ k_{i,j+1/2} \delta_y^-) u_{ij} + O(\Delta x^2, \Delta y^2)$$

- **We could use another form as**

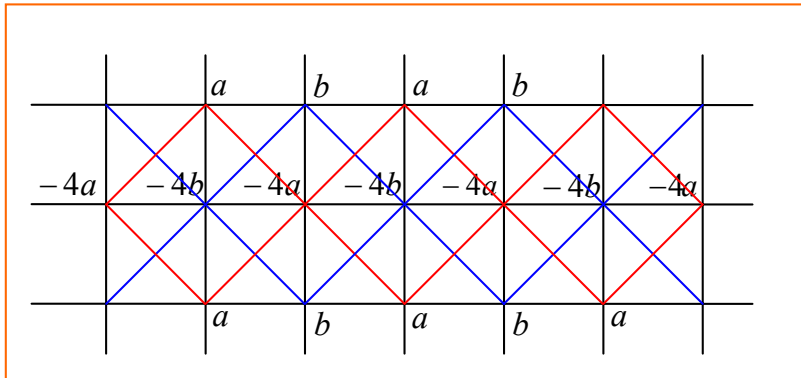
$$\nabla^2 = \frac{1}{\Delta x^2} (\mu_y \delta_x)^2 + \frac{1}{\Delta y^2} (\mu_x \delta_y)^2$$

Multi-dimensional FDM Formulas

- Which gives the following formula for $\Delta x = \Delta y$:

$$\nabla^2 u_{ij} = \frac{1}{4\Delta x^2} (u_{i+1,j+1} + u_{i+1,j-1} + u_{i-1,j-1} + u_{i-1,j+1} - 4u_{i,j})$$

- This formula is not suitable as it decouples the odd-numbered points from the even-numbered points. This means that for a situation where the solution oscillates between values a and b , the results satisfy the Laplace equation while it is not right.



Multi-dimensional FDM Formulas

- We can combine the last two methods to obtain a family of nine-point schemes as ($\Delta x = \Delta y$)

$$\nabla^2 u_{ij} = \frac{1}{\Delta x^2} \left[(\delta_x^2 + \delta_y^2) + \frac{b}{2} \delta_x^2 \delta_y^2 \right] = \Delta u_{ij} + \frac{\Delta x^2}{12} (u_{xxxx} + u_{yyyy} + 6b u_{xxyy})$$

- For $b = 2/3$, this is equivalent to a **Galerkin finite element discretization of the Laplace operator using bilinear quadrilateral elements on the same mesh.**

Mixed Derivatives

- The simplest 2nd order central formula is

$$\begin{aligned}u_{xy} &= \frac{1}{\Delta x \Delta y} \mu_x \delta_x \left[\left(1 + \frac{\delta_x^2}{6} \right) + O(\Delta x^4) \right] \mu_y \delta_y \left[\left(1 + \frac{\delta_y^2}{6} \right) + O(\Delta y^4) \right] u_{ij} \\ &= \frac{1}{\Delta x \Delta y} (\mu_x \delta_x \mu_y \delta_y) u_{ij} + O(\Delta x^2, \Delta y^2) \\ &= \frac{1}{4 \Delta x \Delta y} (u_{i+1, j+1} - u_{i+1, j-1} - u_{i-1, j+1} + u_{i-1, j-1}) + O(\Delta x^2, \Delta y^2)\end{aligned}$$

- A first order formula in both x and y is

$$(u_{xy})_{ij} = \frac{1}{\Delta x \Delta y} \delta_x^+ \delta_y^+ u_{ij} + O(\Delta x, \Delta y) = \frac{1}{\Delta x \Delta y} (u_{i+1, j+1} - u_{i+1, j} - u_{i, j+1} + u_{ij}) + O(\Delta x, \Delta y)$$

- We should note that the above formula is 2nd order at $(i + 1/2, j + 1/2)$.
- The following formula is the most general 2nd order mixed derivative formula: ($a + b = 1$)

$$(u_{xy})_{ij} = \frac{1}{2 \Delta x \Delta y} \delta_x \delta_y (a u_{i+1/2, j+1/2} + a u_{i-1/2, j-1/2} + b u_{i+1/2, j-1/2} + b u_{i-1/2, j+1/2}) + O(\Delta x^2, \Delta y^2)$$

NON-UNIFORM GRIDS

- In a general mesh, Δx may vary. So, we should be careful of the order of the approximations:

$$\begin{aligned}(u_x)_i &= \frac{u_{i+1} - u_i}{\Delta x_{i+1}} - \frac{\Delta x_{i+1}}{2} u_{xx} && \text{(first order)} \\ (u_x)_i &= \frac{u_i - u_{i-1}}{\Delta x_i} + \frac{\Delta x_i}{2} u_{xx} && \text{(first order)}\end{aligned}$$

- Combining these formulas, one gets:

$$(u_x)_i = \frac{1}{\Delta x_i + \Delta x_{i+1}} \left[\frac{\Delta x_i}{\Delta x_{i+1}} (u_{i+1} - u_i) + \frac{\Delta x_{i+1}}{\Delta x_i} (u_i - u_{i-1}) \right] - \frac{\Delta x_i \Delta x_{i+1}}{6} u_{xxx} \quad (2nd \text{ order})$$

- A 2nd order formula for the 2nd derivative is

$$\begin{aligned}(u_{xx})_i &= \left(\frac{u_{i+1} - u_i}{\Delta x_{i+1}} - \frac{u_i - u_{i-1}}{\Delta x_i} \right) \frac{2}{\Delta x_{i+1} + \Delta x_i} \\ &\quad + \frac{1}{3} (\Delta x_{i+1} - \Delta x_i) u_{xxx} - \frac{\Delta x_{i+1}^3 + \Delta x_i^3}{12(\Delta x_{i+1} + \Delta x_i)} u_{xxxx}\end{aligned}$$

NON-UNIFORM GRIDS

- Note that if the mesh size varies abruptly, for example if $\Delta x_{i+1} \approx 2\Delta x_i$, then the formula will be only 1st order accurate.
- It is a general property of finite difference method on non-uniform meshes that if the mesh size does not vary smoothly, a loss of accuracy is unavoidable.